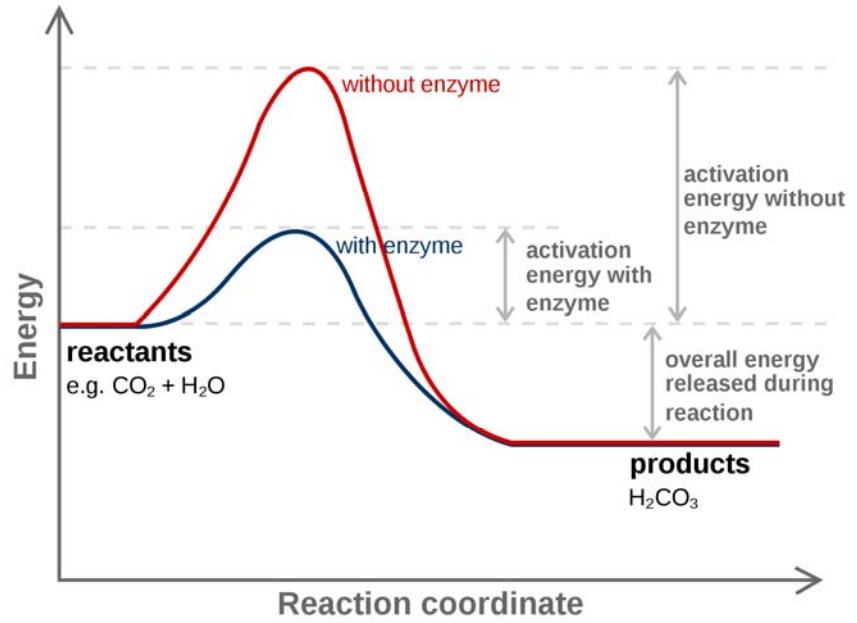
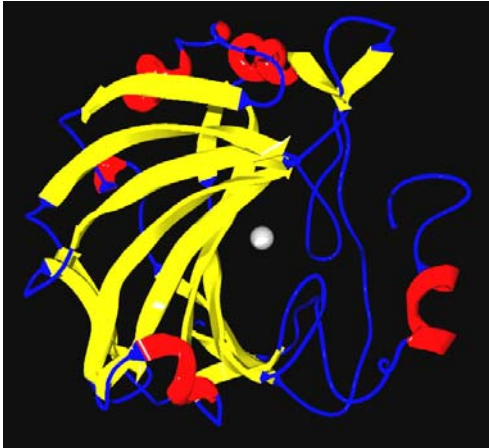


Enzymes are proteins that catalyze (*i.e.*, increase the rate of) chemical reaction

Carbonic anhydrase



(in tissues with high CO_2 concentration)



(in lungs with low CO_2 concentration)

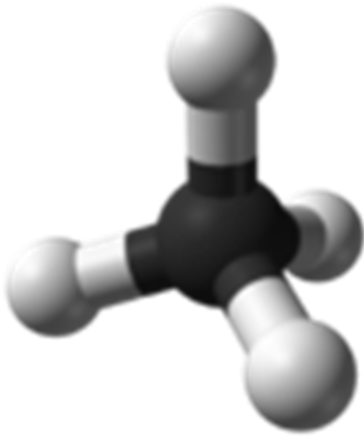
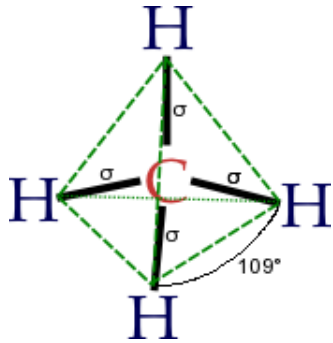
In humans it is working nicely

Can we understand how does it work?

Increasing the rate of chemical reaction is the dream of industry

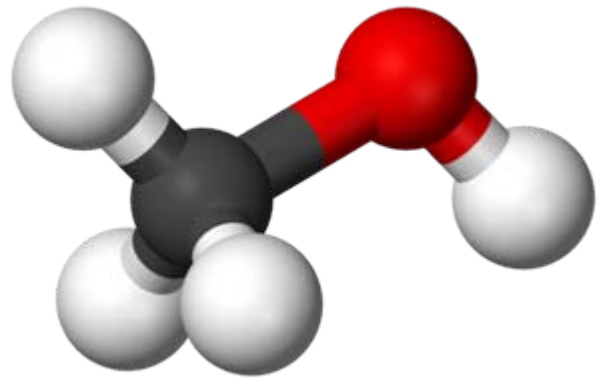
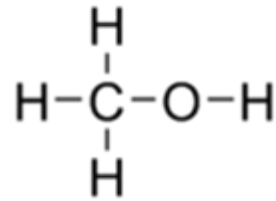
Oxidation of methane to methanol

methane



Gas
Explosive
Huge amount

methanol



Liquid
Easy to storage
fuel

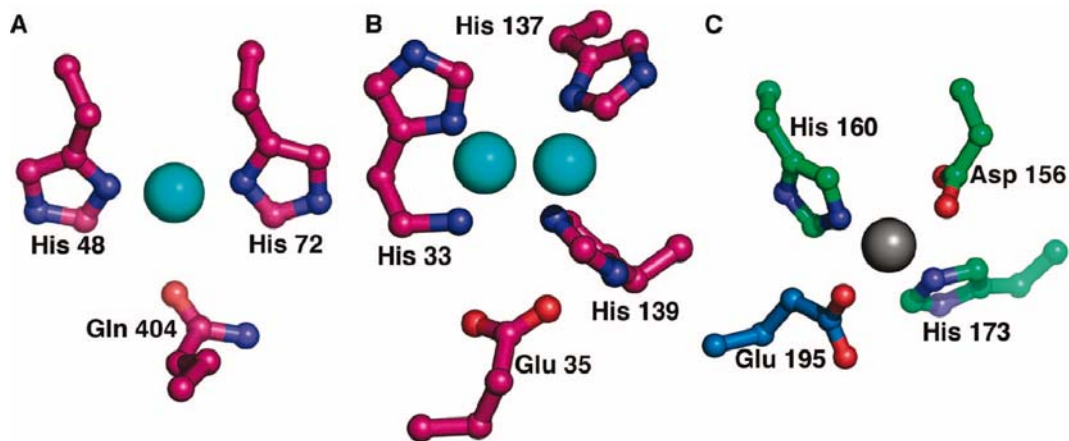
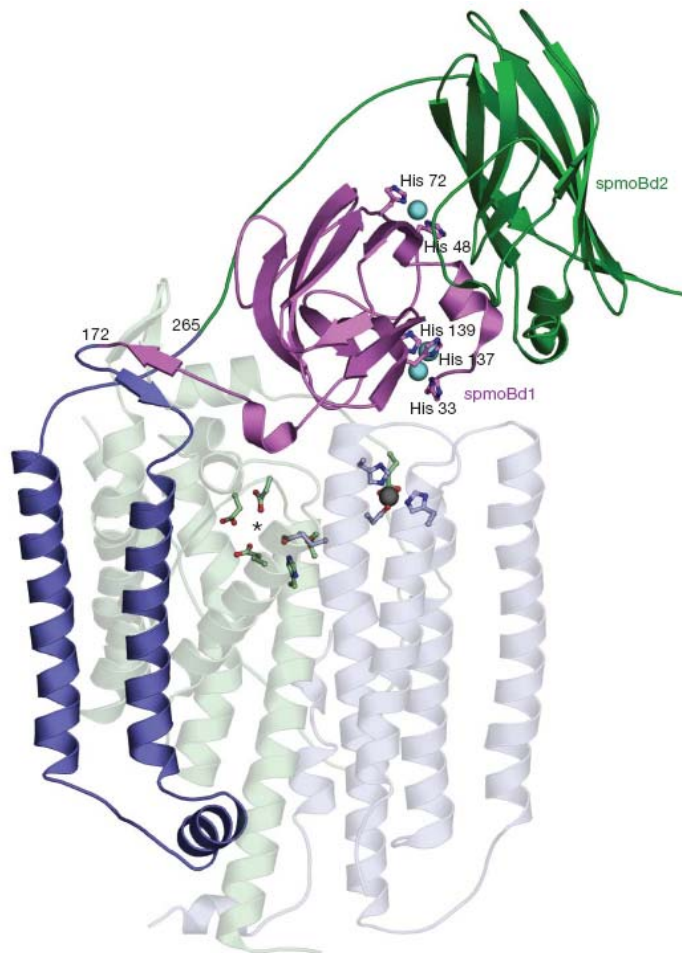
but

The methane oxidation is a difficult and expensive reaction because the methane is the most inert hydrocarbon (C-H bond is one of the most strong)

however

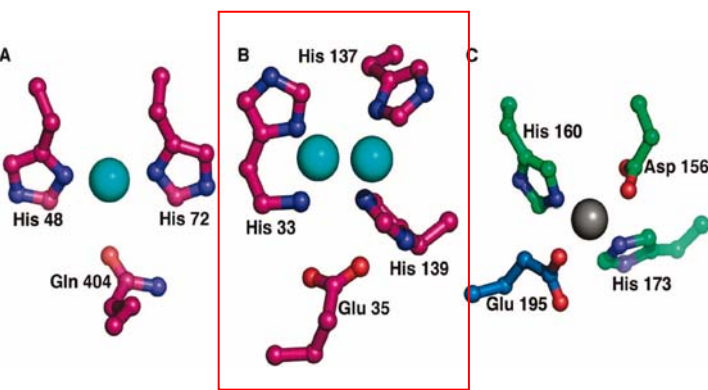
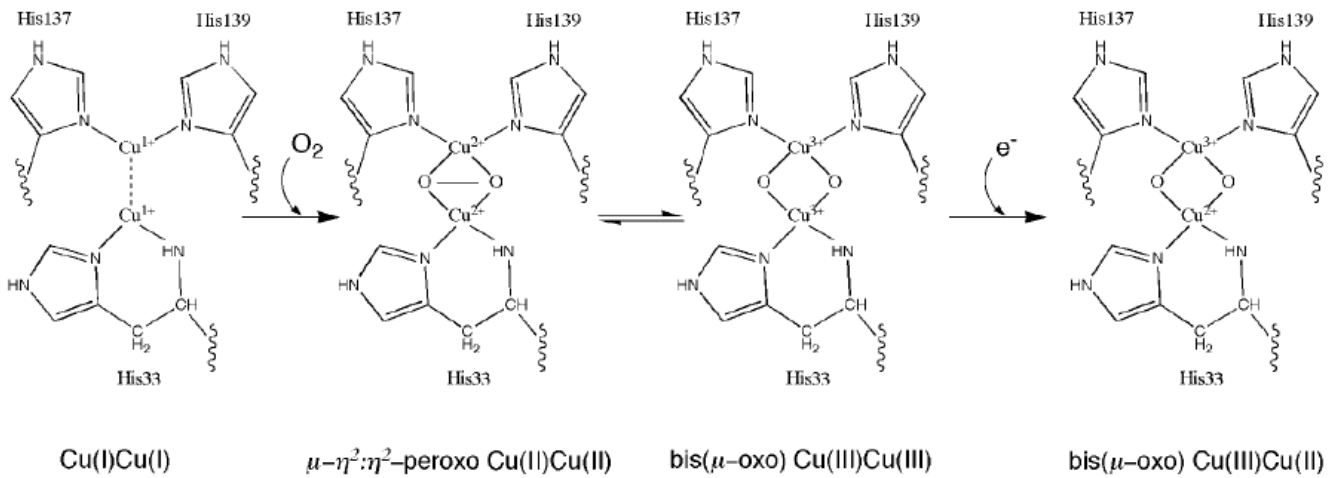
Methanotrophs bacteria use methane monooxygenase (MMO) enzymes to convert methane to methanol at ambient conditions !!!

monooxygenase (MMO)

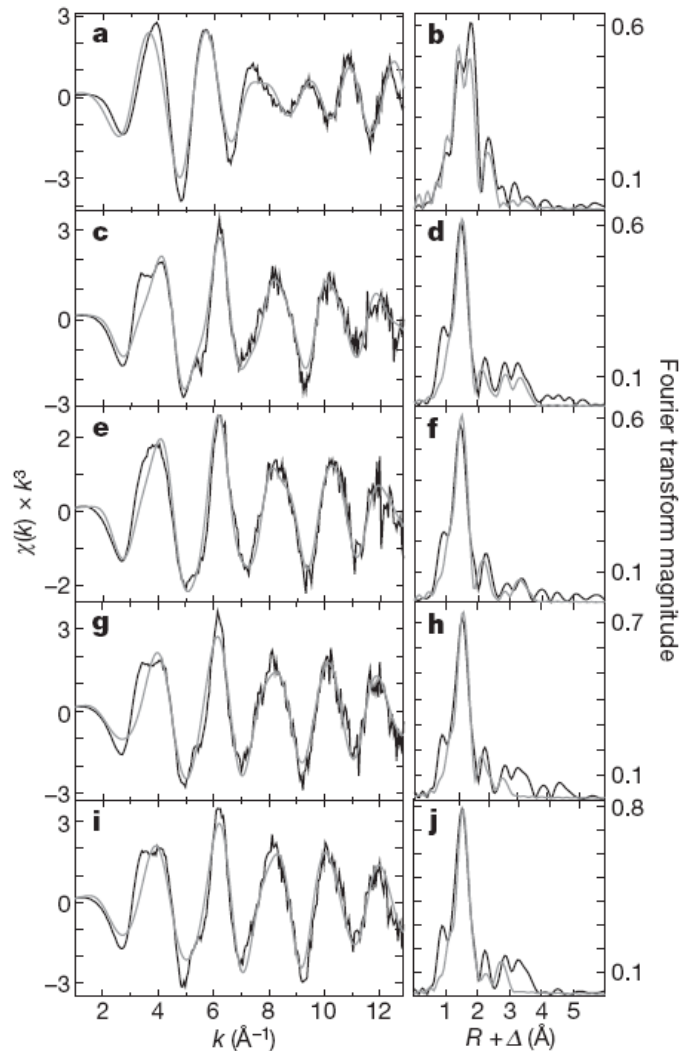


MMO metal centers. (A) Mononuclear copper center. (B) Dinuclear copper center. (C) Zinc center. Oxygen atoms are colored red; nitrogen atoms are colored blue; Zinc atom in grey

Oxidation at the dinuclear copper center



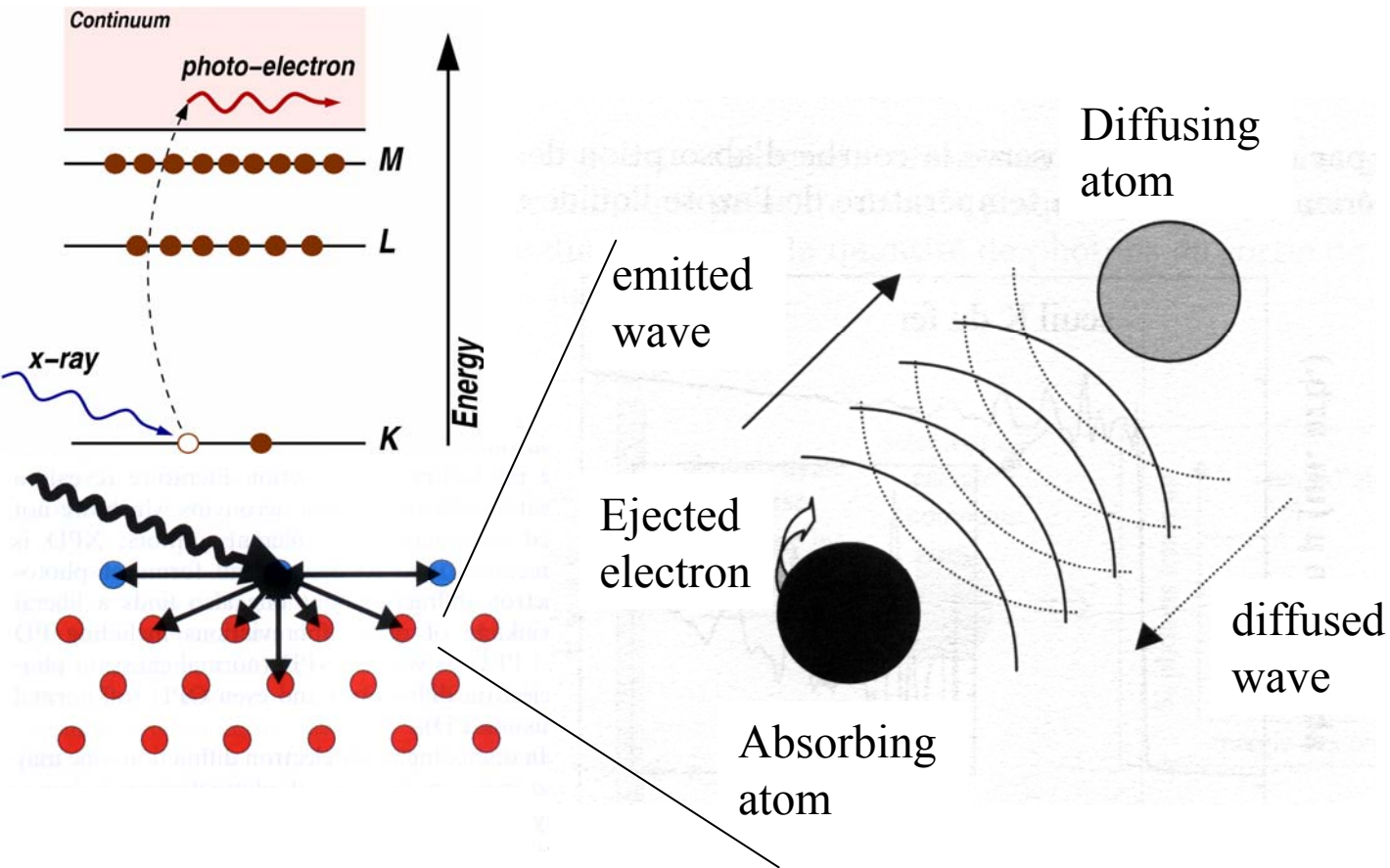
Copper edge EXAFS data and simulations for MMO



EXAFS: Extended X-rays absorption fine structure

Principle:

- 1) photon in - electron out
- 2) The emitted electron (described by a wave) is diffused by the neighboring atoms
- 3) Interference between emitted and diffused wave



The interference between emitted and diffused wave is constructive or destructive depending on:

- 1) The distance between the absorbing and diffusing atoms -> sensibility to the crystallographic structure
- 2) The reflection coefficient of the diffusing atom -> sensibility to the chemical environment

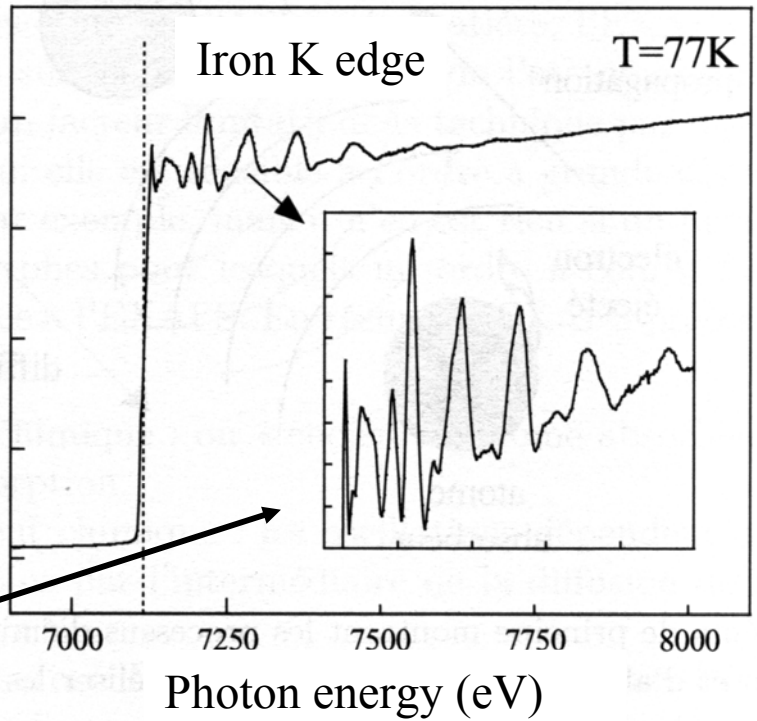
- 3) The wavelength of the emitted electron ->

$$E_{\text{kin}} = h\nu - E_{\text{edge}} = \frac{\hbar^2 k^2}{2m} = \frac{h^2}{(2m\lambda^2)}$$

$$\lambda = h/(2m(h\nu - E_K)^{1/2})$$



absorption (arb. un.)



Interference conditions depend on the in-coming photons

EXAFS expression

$$\chi(k) = (\mu(k) - \mu_0(k))/\mu_0(k) = \frac{1}{k} \sum_i^{\text{atoms}} S_0^2 N_i(\underline{\epsilon}) / R_i^2 A_i(k) \exp(-2R_i/\lambda(k)) \exp(-2k^2\sigma_i^2) \sin(2kR_i + 2\delta_i + \phi_i)$$

($\mu_0(k)$ is the atomic background)

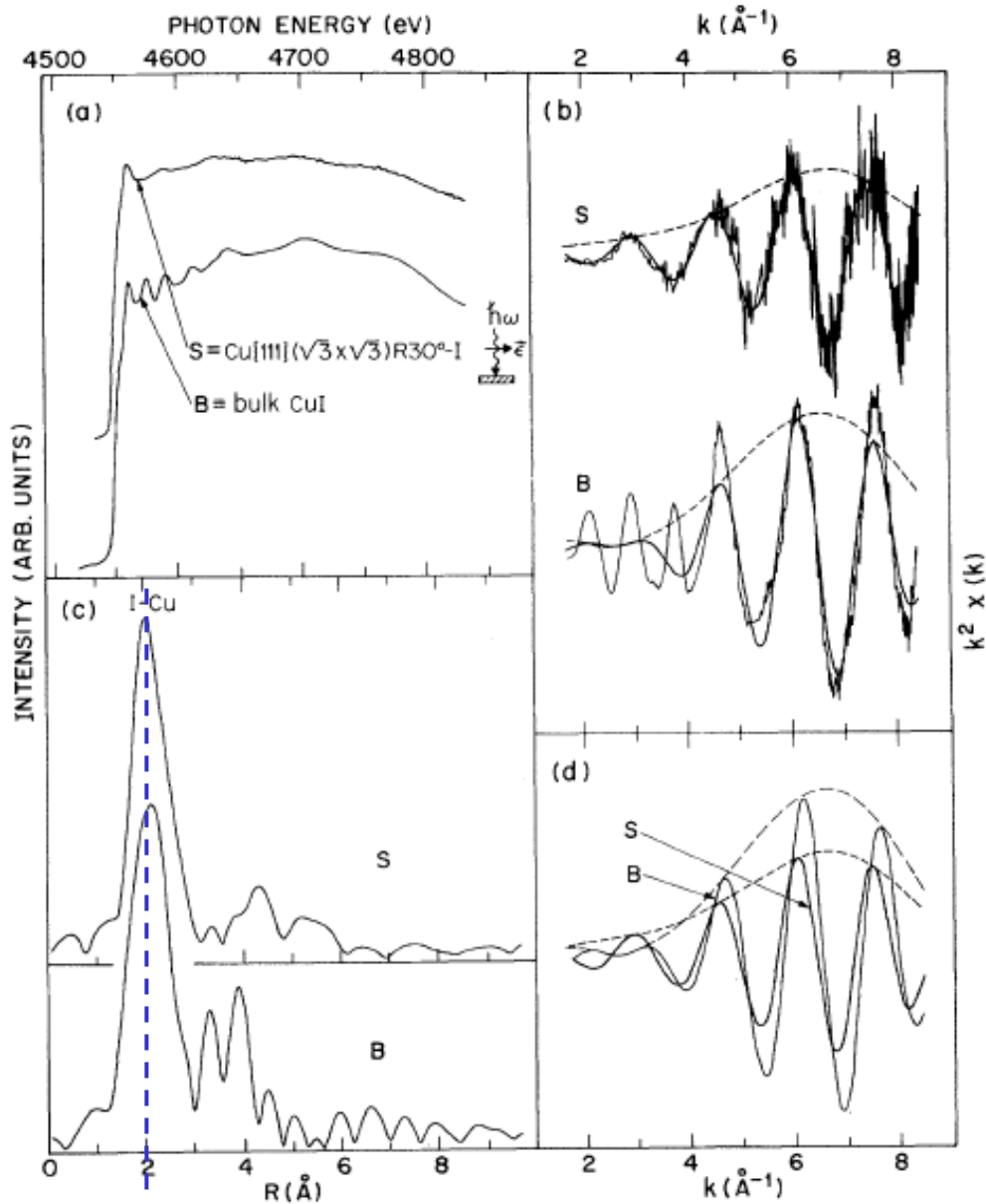
- | | | | |
|------------------|--|--|--------------|
| Amplitude factor | Attenuation factor: reduced electronic mean free path | Debye-Waller term: thermal vibrations and crystallographic disorder | Phase factor |
|------------------|--|--|--------------|

- kR_i -> phase due to the distance between absorbing and diffusing atom
- δ_i -> phase due to the propagation in the potential of the absorbing atom
- ϕ_i -> phase due to the propagation in the potential of the diffusing atom

$k \chi(k)$ is a summation over sinusoid functions

The Fourier transform of $k \chi(k)$ gives information on the phase factor or on the atomic distance

Iodine on Cu(111) vs bulk CuI alloy



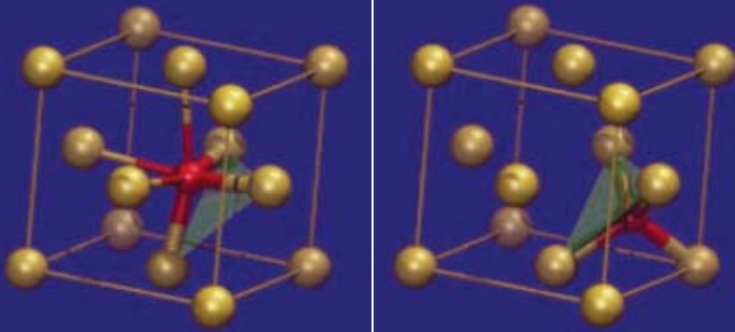
a) and b) EXAFS (at the I edge) and SEXAFS data for bulk CuI (B) and I adsorbed on Cu(111) (S); c) Fourier Transform of data in b). The double peaks for the bulk CuI are due to the I-I bonds; d) Retransformed data from c)

The Cu-I distance for the adsorbed structure is slightly shorter than in the bulk alloy

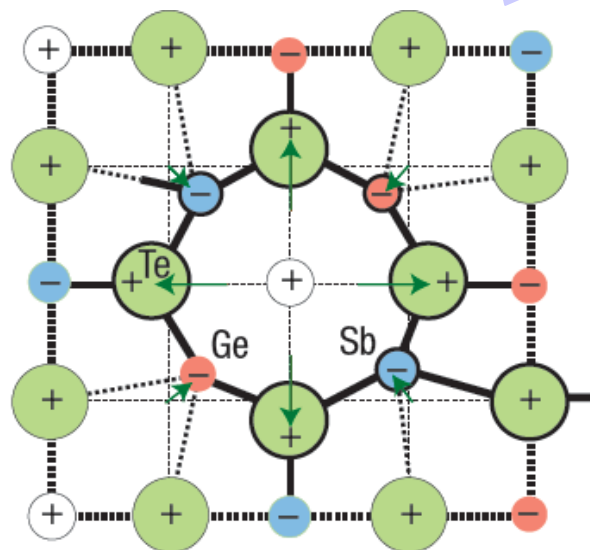
Understanding the phase-change mechanism of rewritable optical media

A. V. Kolobov, *et al.* Nat. Mater. **3**, 703 (2004).

Present interest in chalcogenide glasses is driven by the ability of a particular composition ($\text{Ge}_2\text{Sb}_2\text{Te}_5$, or GST) to be **repeatedly switched between crystalline (*c*) and amorphous (*a*) states** by application of light or electrical pulses of suitable intensities and durations. Differences in the properties of the *a* and *c* materials (e.g., reflectivity) allows for device applications.

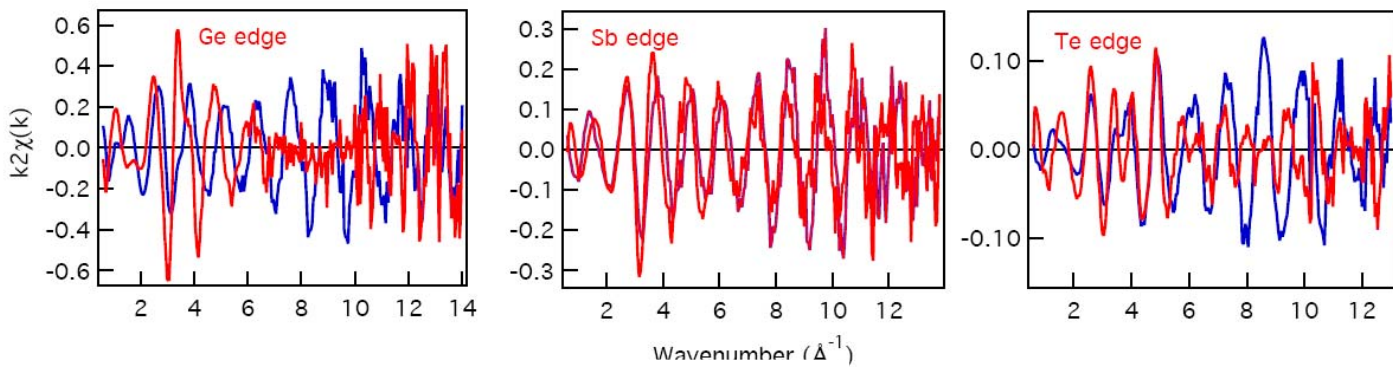


Fragments of the local structure of GST around Ge atoms in the crystalline (left) and amorphous (right) states. Stronger covalent bonds are shown as thicker lines whereas weak interblock bonds are shown as thinner lines.

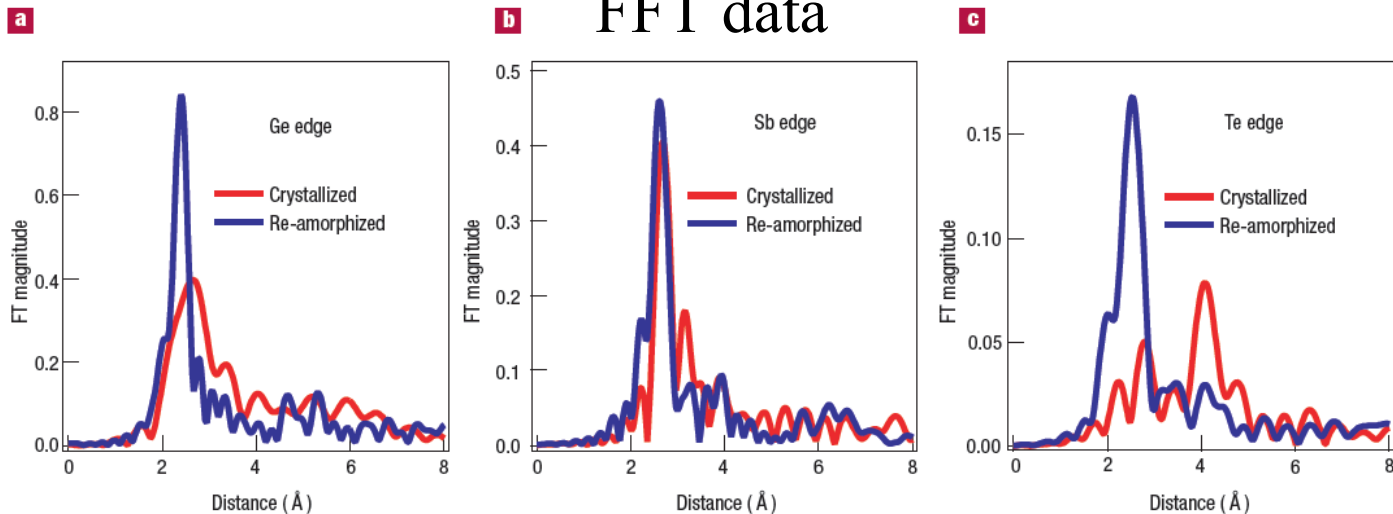


The crystal structure of laser-amorphized GST. A schematic twodimensional image of the lattice distortion of the rocksalt structure due to charge redistribution between the constituent elements; atoms that form the building block of the GST structure are shown using thick lines. The arrows indicate displacements of atoms from the ideal rocksalt positions.

Raw data



FFT data



Spectra measured at the K-edges of: **a**, Ge, **b**, Sb and **c**, Te. On amorphization the **bonds become shorter** (as shown by shifts in the peak positions) and stronger, that is, **more locally ordered** (as shown by increases in the peak amplitudes and concurrent decreases in the peak widths). The peak positions are shifted from the actual interatomic distances towards lower r because of the photoelectron phase shift $\delta(k)$ in the phase factor of the EXAFS oscillations. Additionally, contributions from different nearest neighbors interfere producing extra features in the FTs.

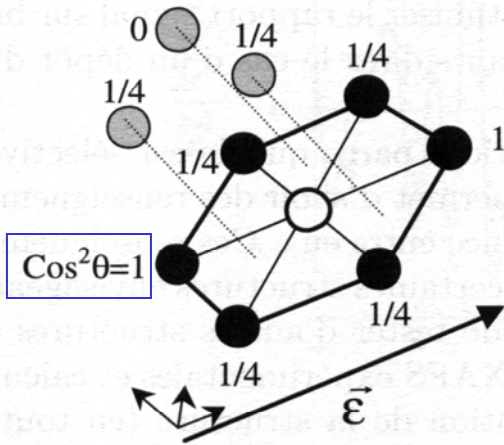
Table 1 Comparison of the GST bond lengths determined from EXAFS and XRD analysis for the crystalline and laser-amorphized states.

| Bond | Bond length (Å) | |
|-------------|-------------------------------|------------------|
| | From EXAFS | From XRD |
| | Crystallized state | |
| Ge–Te | 2.83 ± 0.01 | $3.0(1) \pm 0.3$ |
| Sb–Te | 2.91 ± 0.01 | $3.0(1) \pm 0.3$ |
| Te–Te (2nd) | 4.26 ± 0.01 | $4.2(6) \pm 0.2$ |
| | Laser-amorphized state | |
| Ge–Te | 2.61 ± 0.01 | 2.61* |
| Sb–Te | 2.85 ± 0.01 | |

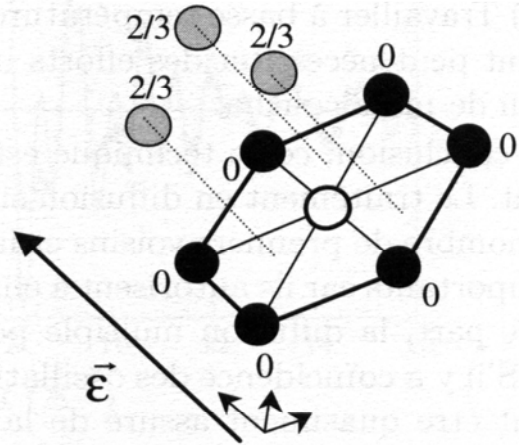
Polarization counts

$$N_i(\underline{\epsilon}) = 3 (\underline{R}_i \cdot \underline{\epsilon})^2 / |\underline{R}_i|^2 |\underline{\epsilon}|^2 = 3 \cos^2 \theta$$

In-plane polarization



Out-of-plane polarization



- Absorbing atom
- 1st Diffusing atom in the plane
- 1st Diffusing atom out of plane

Hcp: hexagonal closed packed

2 overlapping "simple hexagonal lattices"

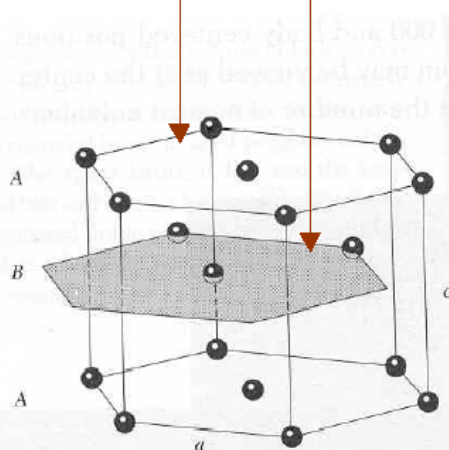


Figure 22 The hexagonal close-packed structure. The atom positions in this structure do not constitute a space lattice. The space lattice is simple hexagonal with a basis of two identical atoms associated with each lattice point. The lattice parameters a and c are indicated, where a is in the basal plane and c is the magnitude of the axis a_3 of Fig. 14.

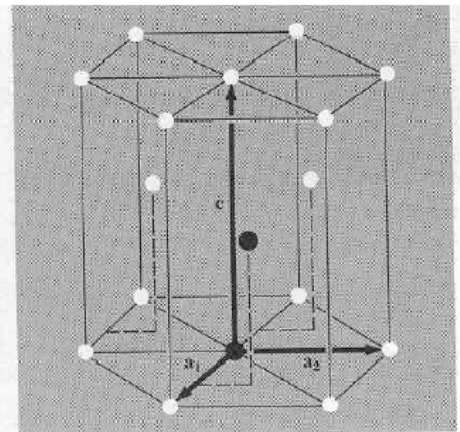
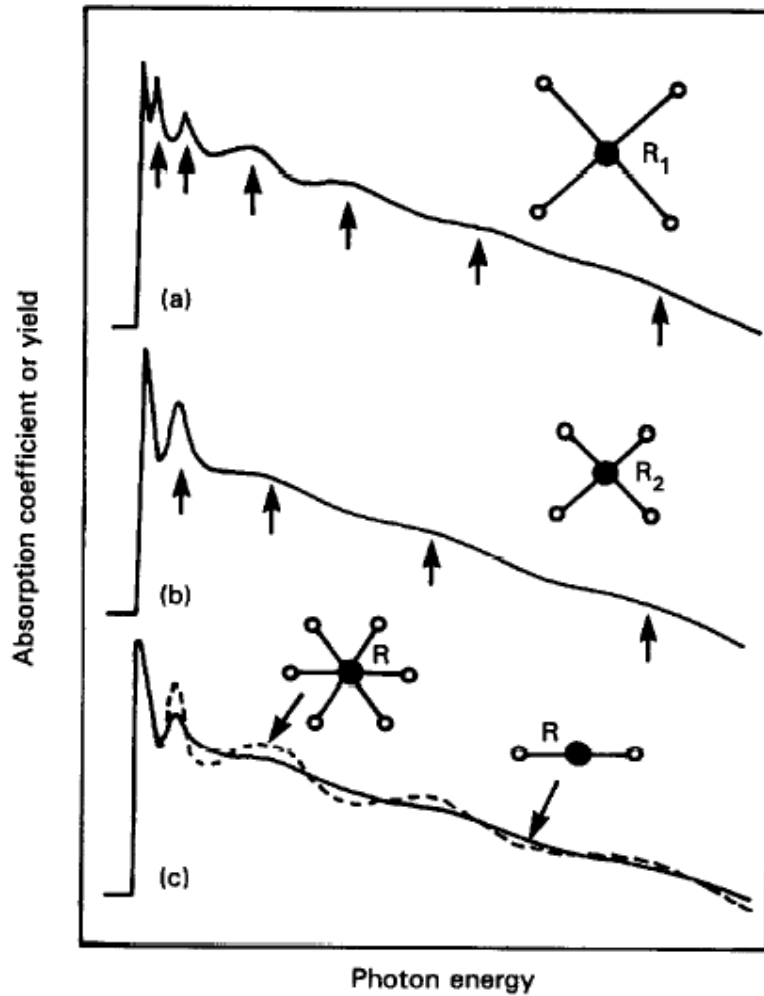


Figure 23 The primitive cell has $a_1 = a_2$, with an included angle of 120° . The c axis (or a_3) is normal to the plane of a_1 and a_2 . The ideal hcp structure has $c = 1.633 a$. The two atoms of one basis are shown as solid circles. One atom of the basis is at the origin; the other atom is at $\frac{2}{3}\mathbf{a}_1 + \frac{1}{3}\mathbf{a}_2 + \frac{1}{2}\mathbf{a}_3$.

Oscillations in absorption due to *interference* between *outgoing and scattered photoelectrons* by neighbors at r

- Higher coordination \rightarrow higher amplitude
- Shorter r \rightarrow increased period



Convention:
 EXAFS \rightarrow about 100 eV above absorption edge
 XANES \rightarrow around the absorption edge

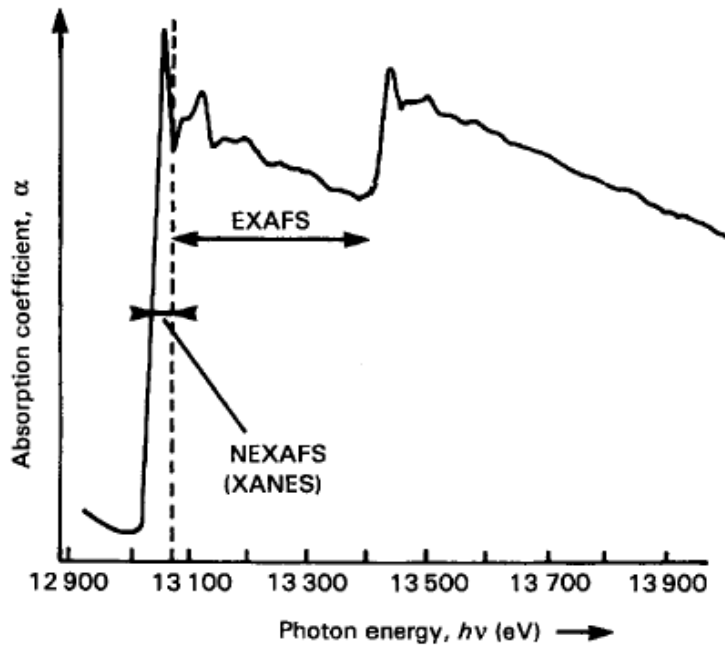
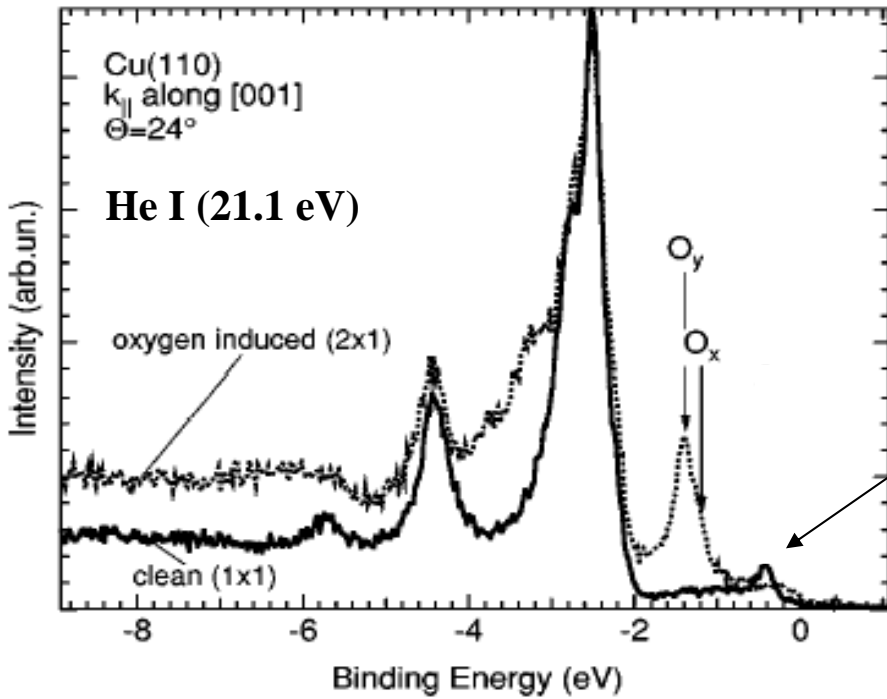


Figure 8.24. A typical X-ray absorption spectrum of a condensed medium. The sample is $\text{BaPb}_{1-x}\text{Bi}_x\text{O}_3$, and the figure shows the L_{III} absorption edge of Pb, with the NEXAFS and EXAFS regions indicated. Absorption above around 13400 eV is due to the L_{III} edge of Bi, which is close in energy to the Pb edge

Cu(110) (2x1) O



UPS spectrum

K. Stahrenberg, *et al.* Phys. Rev. B **61**, 3043 (2000).

Cu (110)
Surface
state

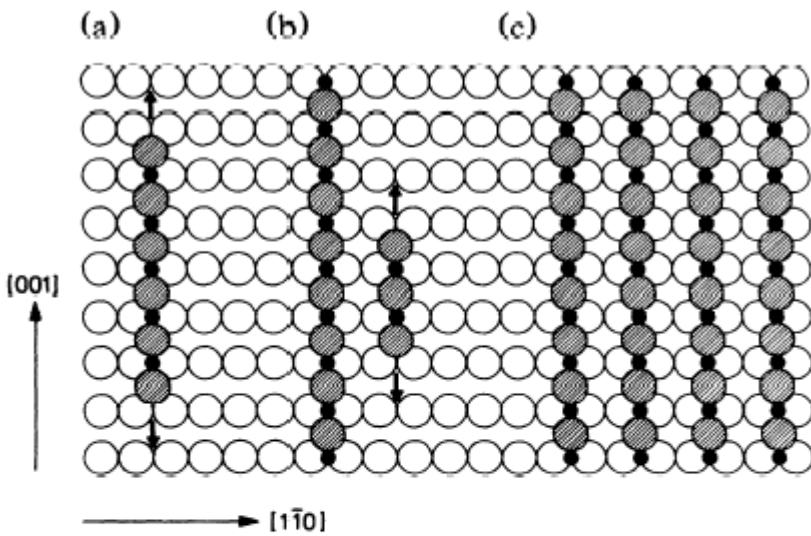


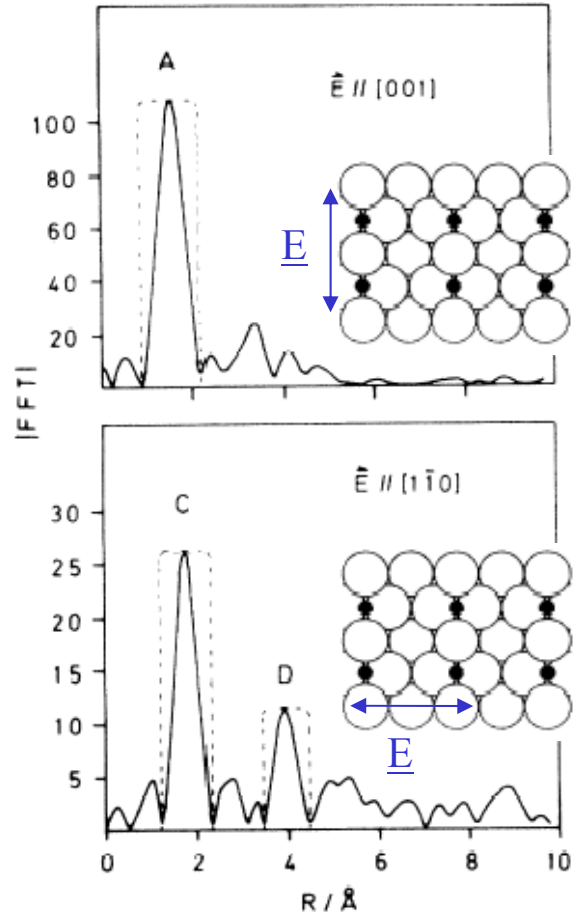
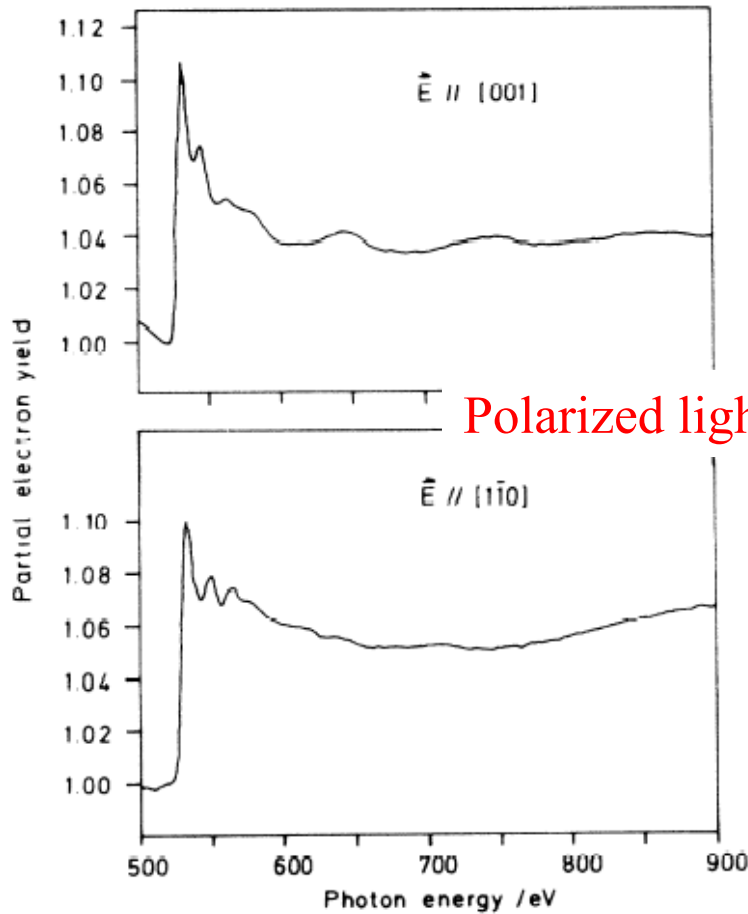
FIG. 4. Atomistic model of the different stages of (2x1)O formation. (a) Single string of Cu-O adatoms along [001] ("added row"); arrows indicate preferential growth direction. (b) Growth of a single-row (2x1)O island along [110]; nucleation of a neighbored added row. (c) Two-dimensional island of (2x1)O added-row phase, the structure being equivalent to the "missing-row" structure. Filled circles: O atoms; shaded circles: added-row Cu atoms on top of the substrate atoms (open circles).



100 Å

STM image of nuclei in different growth phases (step edges are marked by arrows)

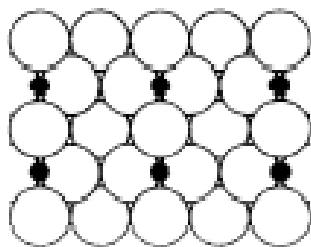
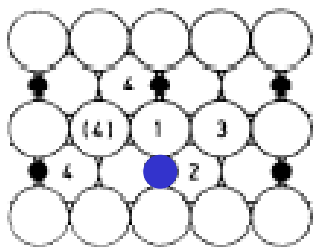
Oxygen-K edges SEXAFS spectra for the Cu(110) (2x1) O at 100 K



| Azimuth | θ /degrees | Cu neighbor | No reconstruction | Missing row |
|---------|-------------------|-------------|-------------------|-------------|
| [001] | 90 | 1 | 1 | 1 |
| [1-10] | 90 | 2 | 0.355 | 0.355 |
| | | 3 | 0.458 | ... |
| | | 4 | 0.354 | 0.354 |

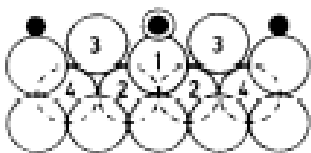
Buckled row

Missing row



→ [1-10]

No peak due to the third neighboring is observed for $E // [1\bar{1}0]$ → missing row structure



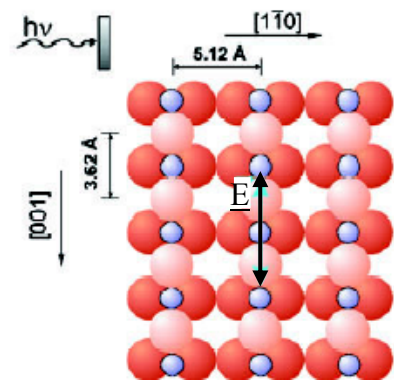
AXAFS: Atomic EXAFS

SEXAFS -> backscattering of the photoelectron at neighboring atoms
 AXAFS -> scattering at the charge densities placed between the atoms.

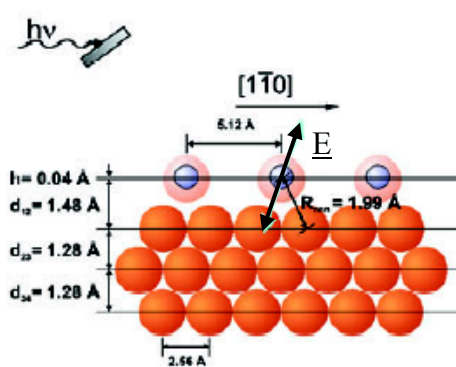
H. Wende *et al.* Rep. Prog. Phys. **67**, 2105 (2004)

Cu(110) (2x1) O

Top view:
normal x-ray
incidence



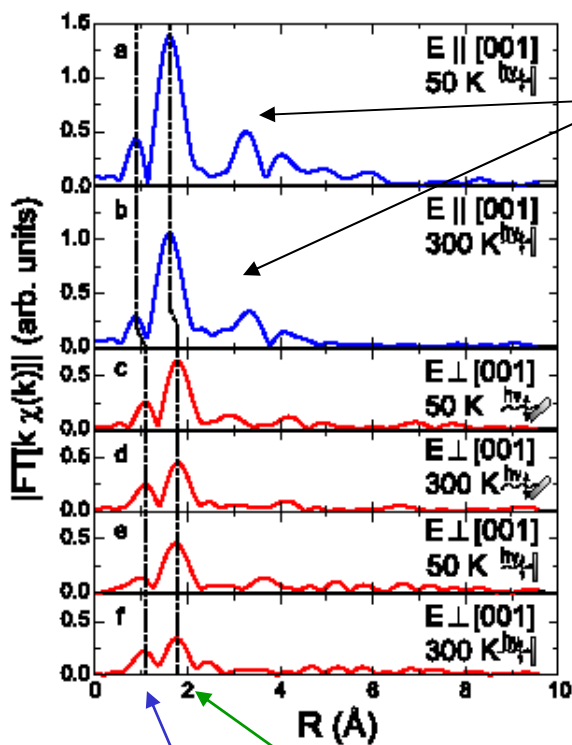
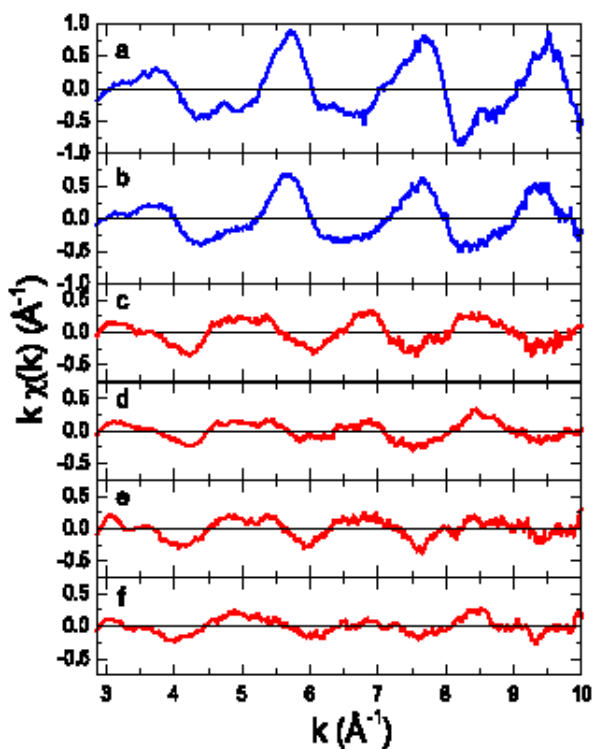
Side view:
grazing x-ray
incidence



Small shaded circles: oxygen atoms, large light circles: Cu atoms of the 1st layer, large dark circles: Cu atoms of the second and following layers. The E -vector for different x-ray incidences is given schematically.

Sensitivity to the O - first layer Cu bond

Sensitivity to the O - second layer Cu bond



Intensity reduced at 300 K in comparison to 50 K due to the Debye-Waller term

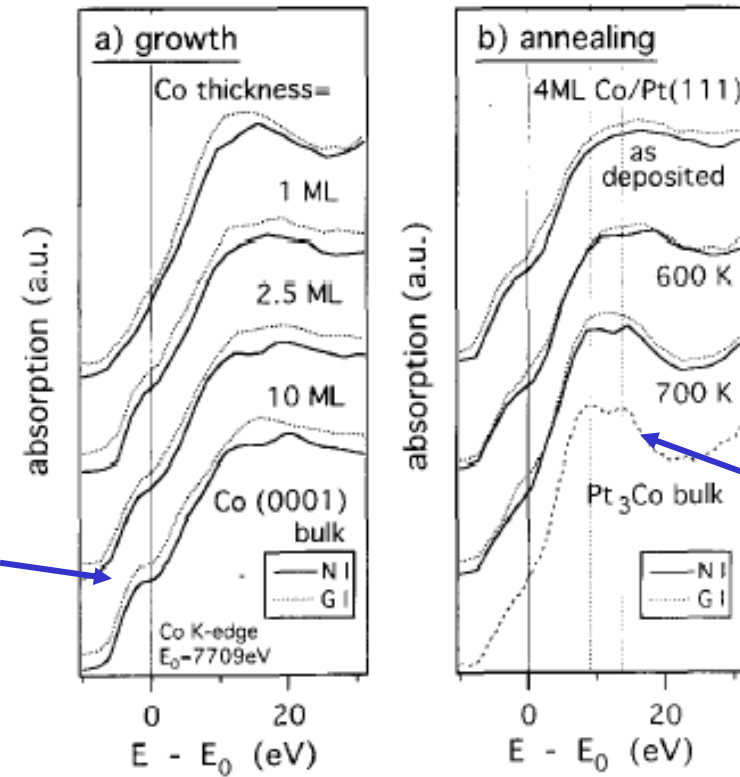
SEXAFS nearest neighbor $1.81 \pm 0.03 \text{ \AA}$
 SEXAFS next nearest neighbor $1.99 \pm 0.03 \text{ \AA}$
 AXAFS nearest neighbor $0.95 \pm 0.05 \text{ \AA}$
 AXAFS next nearest neighbor $1.15 \pm 0.05 \text{ \AA}$

AXAFS

SEXAFS

XANES: X-ray Absorption Near Edge Structure

The absorption edge shape is representative of the film thickness and chemical composition



Co/Pt(111)

J. Thiele *et al.* Surf. Sci. **384**, 120 (1997)

SiO₂/Ni₈₁Fe₁₉(80 Å)/Co(20 Å)/
Al(20 Å+plasma oxidation for x
seconds)/ Ni₁₈Fe₁₉(100 Å)/Cu(30 Å)

Satellite peak due to the CoO

