1) The "Ket" and the associated Dirac or usual vector notations are:

•
$$|H\rangle = \begin{pmatrix} 1\\ 0 \end{pmatrix}$$
 and $\langle H| = \begin{pmatrix} 1 & 0 \end{pmatrix}$
• $|V\rangle = \begin{pmatrix} 0\\ 1 \end{pmatrix}$ and $\langle H| = \begin{pmatrix} 0 & 1 \end{pmatrix}$
• $\alpha |H\rangle + \beta |V\rangle = \begin{pmatrix} \alpha\\ \beta \end{pmatrix}$ and $\alpha^* \langle H| + \beta^* \langle V| = \begin{pmatrix} \alpha^* & \beta^* \end{pmatrix}$

2) In Dirac notation:

$$\begin{split} & (\gamma^* \langle H | + \delta^* \langle V |) \left(\alpha \left| H \right\rangle + \beta \left| V \right\rangle \right) \\ &= \gamma^* \alpha \left\langle H | H \right\rangle + \gamma^* \beta \left\langle H | V \right\rangle + \delta^* \alpha \left\langle V | H \right\rangle + \delta^* \beta \left\langle V | V \right\rangle \\ &= \gamma^* \alpha + \delta^* \beta \end{split}$$

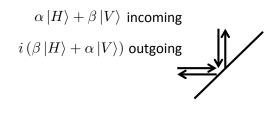
because $\langle H|V\rangle = \langle V|H\rangle = 0$ and $\langle H|H\rangle = \langle V|V\rangle = 1$. The equivalent vector notation is

$$\begin{pmatrix} \gamma^* & \delta^* \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = \gamma^* \alpha + \delta^* \beta.$$

3) We have $R^{\top} = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}$ and $R^{\dagger} = R^{\top,*} = \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix}$. Thus $RR^{\dagger} = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$ $R^{\dagger}R = \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix} \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}.$

Matrices satisfying $MM^{\dagger} = M^{\dagger}M = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$ are called unitary matrices. Let us compute $R(\alpha | H \rangle + \beta | V \rangle)$ in Dirac notation. By linearity of matrix operations,

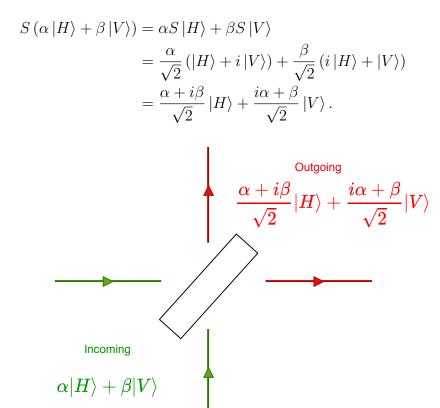
$$R (\alpha |H\rangle + \beta |V\rangle) = \alpha R |H\rangle + \beta R |V\rangle$$
$$= \alpha i |V\rangle + \beta i |H\rangle$$
$$= i (\alpha |V\rangle + \beta |H\rangle).$$



4) We have

$$S |H\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & i \\ i & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix}$$
$$= \frac{1}{\sqrt{2}} \left(|H\rangle + i |V\rangle \right),$$

$$S |V\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & i \\ i & 1 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} i \\ 1 \end{pmatrix}$$
$$= \frac{1}{\sqrt{2}} \left(i |H\rangle + |V\rangle \right),$$



5) The semi-transparent mirror leaves photons in state $S(\alpha |H\rangle + \beta |V\rangle)$. The state is then measured by the detector D which detects photons in state $|V\rangle$. Therefore, the probability of finding a photon in D is the probability of finding a photon in state $|V\rangle$ given that

photons in state $S(\alpha |H\rangle + \beta |V\rangle)$ are produced. By the measurement postulate (which will be formally introduced in Chapter 3 of the lecture note) we have

$$\operatorname{Prob}(D) = \left| \langle V | S(\alpha | H \rangle + \beta | V \rangle) \right|^{2}.$$

From the previous question we have

$$\begin{split} S\left(\alpha\left|H\right\rangle+\beta\left|V\right\rangle\right) &= \frac{\alpha+i\beta}{\sqrt{2}}\left|H\right\rangle + \frac{i\alpha+\beta}{\sqrt{2}}\left|V\right\rangle,\\ \left\langle V\right|S\left(\alpha\left|H\right\rangle+\beta\left|V\right\rangle\right) &= \frac{\alpha+i\beta}{\sqrt{2}}\left\langle V\right|H\right\rangle + \frac{i\alpha+\beta}{\sqrt{2}}\left\langle V\right|V\right\rangle\\ &= \frac{i\alpha+\beta}{\sqrt{2}}. \end{split}$$

So we find

$$Prob(D) = \left|\frac{i\alpha + \beta}{\sqrt{2}}\right|^2 = \frac{1}{2}|i\alpha + \beta|^2 = \frac{1}{2}(\alpha^2 + \beta^2) = \frac{1}{2}.$$

6) The state after S is

$$S|H\rangle = \frac{1}{\sqrt{2}} \left(|H\rangle + i|V\rangle\right)$$

The state after R is

$$RS |H\rangle = \frac{1}{\sqrt{2}} (R |H\rangle + iR |V\rangle)$$
$$= \frac{i}{\sqrt{2}} (|V\rangle + i |H\rangle).$$

The state after the second S is

$$SRS |H\rangle = \frac{i}{\sqrt{2}} \left(S |V\rangle + iS |H\rangle \right)$$
$$= \frac{i}{\sqrt{2}} \left(\frac{i |H\rangle + |V\rangle}{\sqrt{2}} + i \cdot \frac{|H\rangle + i |V\rangle}{\sqrt{2}} \right)$$
$$= -|H\rangle.$$

Thus

$$\operatorname{Prob}(D_1) = |\langle V|H \rangle|^2 = 0$$

$$\operatorname{Prob}(D_2) = |\langle H|H \rangle|^2 = 1.$$

All photons are detected in D_2 ! For "classical balls" we would expect a split between D_1 and D_2 . For example, if S act as half-half splitters we would expect $Prob(D_1) = Prob(D_2) = 1/2$. The quantum behavior is completely different!