
Homework 9: Solution
Quantum Information Processing

Exercise 1 *Bell states*

1) One has to show that $\langle B_{x,y}|B_{x',y'}\rangle = \delta_{x,x'}\delta_{y,y'}$. We show it explicitly for two cases:

$$\begin{aligned}\langle B_{00}|B_{00}\rangle &= \frac{1}{2}(\langle 00| + \langle 11|)(|00\rangle + |11\rangle) \\ &= \frac{1}{2}(\langle 00|00\rangle + \langle 00|11\rangle + \langle 11|00\rangle + \langle 11|11\rangle).\end{aligned}$$

Now we have

$$\begin{aligned}\langle 00|00\rangle &= \langle 0|0\rangle \langle 0|0\rangle = 1, \langle 00|11\rangle = \langle 0|1\rangle \langle 0|1\rangle = 0, \\ \langle 11|00\rangle &= \langle 1|0\rangle \langle 1|0\rangle = 0, \langle 11|11\rangle = \langle 1|1\rangle \langle 1|1\rangle = 1.\end{aligned}$$

Thus we get that $\langle B_{00}|B_{00}\rangle = \frac{1}{2}(1 + 0 + 0 + 1) = 1$. Now let us consider

$$\begin{aligned}\langle B_{00}|B_{01}\rangle &= \frac{1}{2}(\langle 00| + \langle 11|)(|01\rangle + |10\rangle) \\ &= \frac{1}{2}(\langle 00|01\rangle + \langle 00|10\rangle + \langle 11|01\rangle + \langle 11|10\rangle) \\ &= \frac{1}{2}(0 + 0 + 0 + 0) = 0.\end{aligned}$$

2) The proof is by contradiction. Suppose there exists a_1, b_1 and a_2, b_2 such that

$$|B_{00}\rangle = (a_1|0\rangle + b_1|1\rangle) \otimes (a_2|0\rangle + b_2|1\rangle).$$

Then we have

$$\frac{1}{2}(|00\rangle + |11\rangle) = a_1a_2|00\rangle + a_1b_2|01\rangle + b_1a_2|10\rangle + a_2b_2|11\rangle.$$

Comparing the coefficients of the orthonormal basis, one has

$$\frac{1}{2} = a_1a_2, \quad \frac{1}{2} = b_1b_2, \quad a_1b_2 = 0, \quad b_1a_2 = 0.$$

The third equality indicates that either $a_1 = 0$ or $b_2 = 0$ (or both). If $a_1 = 0$ we get a contradiction with the first equation. If on the other hand $b_2 = 0$, we get a contradiction with the second one. Therefore, there does not exist $|\psi_1\rangle$ and $|\psi_2\rangle$ such that $|B_{00}\rangle$ can be written as $|\psi_1\rangle \otimes |\psi_2\rangle$. Therefore, B_{00} is entangled.

3) We have

$$\begin{aligned} |\gamma\rangle \otimes |\gamma\rangle &= (\cos(\gamma)|0\rangle + \sin(\gamma)|1\rangle) \otimes (\cos(\gamma)|0\rangle + \sin(\gamma)|1\rangle) \\ &= \cos^2(\gamma)|00\rangle + \cos(\gamma)\sin(\gamma)|01\rangle + \sin(\gamma)\cos(\gamma)|10\rangle + \sin^2(\gamma)|11\rangle. \end{aligned}$$

Similarly, we have

$$|\gamma_\perp\rangle \otimes |\gamma_\perp\rangle = \sin^2(\gamma)|00\rangle - \cos(\gamma)\sin(\gamma)|01\rangle - \sin(\gamma)\cos(\gamma)|10\rangle + \cos^2(\gamma)|11\rangle.$$

Combining the two terms, we find that

$$\begin{aligned} |\gamma\rangle \otimes |\gamma\rangle + |\gamma_\perp\rangle \otimes |\gamma_\perp\rangle &= (\cos^2(\gamma) + \sin^2(\gamma))|00\rangle + (\sin^2(\gamma) + \cos^2(\gamma))|11\rangle \\ &= |00\rangle + |11\rangle \end{aligned}$$

and thus

$$\frac{1}{\sqrt{2}}(|\gamma\rangle \otimes |\gamma\rangle + |\gamma_\perp\rangle \otimes |\gamma_\perp\rangle) = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) = |B_{00}\rangle.$$

4) From the rule of the tensor product

$$\begin{pmatrix} a \\ b \end{pmatrix} \otimes \begin{pmatrix} c \\ d \end{pmatrix} = \begin{pmatrix} a \begin{pmatrix} c \\ d \end{pmatrix} \\ b \begin{pmatrix} c \\ d \end{pmatrix} \end{pmatrix} = \begin{pmatrix} ac \\ ad \\ bc \\ bd \end{pmatrix},$$

we obtain the basis states as

$$\begin{aligned} |0\rangle \otimes |0\rangle &= \begin{pmatrix} 1 \\ 0 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}, & |0\rangle \otimes |1\rangle &= \begin{pmatrix} 1 \\ 0 \end{pmatrix} \otimes \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix}, \\ |1\rangle \otimes |0\rangle &= \begin{pmatrix} 0 \\ 1 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}, & |1\rangle \otimes |1\rangle &= \begin{pmatrix} 0 \\ 1 \end{pmatrix} \otimes \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}. \end{aligned}$$

Thus, we have

$$\begin{aligned} |B_{00}\rangle &= \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 1 \end{pmatrix}, & |B_{01}\rangle &= \frac{1}{\sqrt{2}}(|01\rangle + |10\rangle) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ 1 \\ 0 \end{pmatrix}, \\ |B_{10}\rangle &= \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ 0 \\ -1 \end{pmatrix}, & |B_{11}\rangle &= \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ -1 \\ 0 \end{pmatrix}. \end{aligned}$$

Exercise 2 *Entanglement by unitary operations*

1) By definition of the tensor product:

$$(H \otimes I) |x\rangle \otimes |y\rangle = H |x\rangle \otimes I |y\rangle = H |x\rangle \otimes |y\rangle.$$

Also, one can use that $H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$ to show that always

$$H |x\rangle = \frac{1}{\sqrt{2}}(|0\rangle + (-1)^x |1\rangle).$$

Thus,

$$(H \otimes I) |x\rangle \otimes |y\rangle = \frac{1}{\sqrt{2}}(|0\rangle \otimes |y\rangle + (-1)^x |1\rangle \otimes |y\rangle).$$

Now we apply CNOT. By linearity, we can apply it to each term separately. Thus,

$$\begin{aligned} (\text{CNOT})(H \otimes I) |x\rangle \otimes |y\rangle &= \frac{1}{\sqrt{2}}((\text{CNOT}) |0\rangle \otimes |y\rangle + (-1)^x (\text{CNOT}) |1\rangle \otimes |y\rangle) \\ &= \frac{1}{\sqrt{2}}(|0\rangle \otimes |y\rangle + (-1)^x |1\rangle \otimes |y \oplus 1\rangle) \\ &= |B_{xy}\rangle. \end{aligned}$$

2) Let us first start with $H \otimes I$. We use the rule

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} \otimes \begin{pmatrix} e & f \\ g & h \end{pmatrix} = \begin{pmatrix} ae & af & be & bf \\ ag & ah & bg & bh \\ ce & cf & de & df \\ cg & ch & dg & dh \end{pmatrix},$$

Thus we have

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \otimes \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 & 1 & 0 \\ 0 & 1 & 0 & 1 \\ 1 & 0 & -1 & 0 \\ 0 & 1 & 0 & -1 \end{pmatrix}.$$

For CNOT, we use the definition:

$$(\text{CNOT}) |x\rangle \otimes |y\rangle = |x\rangle \otimes |y \oplus x\rangle,$$

which implies that the matrix elements are

$$\langle x'y' | \text{CNOT} |xy\rangle = \langle x', y' | x, y \oplus x\rangle = \langle x' | x\rangle \langle y' | y \oplus x\rangle = \delta_{xx'} \delta_{y \oplus x, y'}.$$

We obtain the following table with columns xy and rows $x'y'$:

	00	01	10	11
00	1	0	0	0
01	0	1	0	0
10	0	0	0	1
11	0	0	1	0

For the matrix product $(\text{CNOT})(H \otimes I)$, we find that

$$\begin{aligned} (\text{CNOT})H \otimes I &= \frac{1}{\sqrt{2}} \begin{pmatrix} I & 0 \\ 0 & X \end{pmatrix} \begin{pmatrix} I & I \\ I & -I \end{pmatrix} \\ &= \frac{1}{\sqrt{2}} \begin{pmatrix} I & I \\ X & -X \end{pmatrix}, \end{aligned}$$

where $X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$. Thus,

$$(\text{CNOT})(H \otimes I) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 & 1 & 0 \\ 0 & 1 & 0 & 1 \\ 0 & 1 & 0 & -1 \\ 1 & 0 & -1 & 0 \end{pmatrix}.$$

One can check that for example $|B_{00}\rangle = (\text{CNOT})(H \otimes I)|0\rangle \otimes |0\rangle$. Finally to check the unitarity, we have to check that $UU^\dagger = U^\dagger U = I$ for $U = H \otimes I$, CNOT and $(\text{CNOT})(H \otimes I)$. We leave this to the reader.

3) Let $U = (\text{CNOT})(H \otimes I)$. We have

$$|B_{xy}\rangle = U|x\rangle \otimes |y\rangle, \langle B_{x'y'}| = \langle x'| \otimes \langle y'| U^\dagger.$$

Thus,

$$\begin{aligned} \langle B_{x'y'}|B_{xy}\rangle &= \langle x'| \otimes \langle y'| U^\dagger U |x\rangle \otimes |y\rangle \\ &= \langle x'| \otimes \langle y'| I |x\rangle \otimes |y\rangle \\ &= \langle x'|x\rangle \langle y'|y\rangle = \delta_{xx'} \delta_{yy'}. \end{aligned}$$

Exercise 3 Tsirelson bound

1) Since the eigenvalues of A are ± 1 , those of A^2 are both $+1$. The eigenvectors are the same. Thus we have $A^2 = |\alpha\rangle\langle\alpha| + |\alpha_\perp\rangle\langle\alpha_\perp| = I_A$ and similarly for the other matrices. The same result can also be obtained by expanding the product A^2 in Dirac notation or in matrix component form...

Now expanding \mathcal{B}^2 we get the squared terms of the type $(A \otimes B)^2 = (A \otimes B)(A \otimes B) = A^2 \otimes B^2 = I_A \otimes I_B$. This will yield the term $4I_A \otimes I_B$.

Then for the cross terms using $(M_1 \otimes M_2)(N_1 \otimes N_2) = M_1 N_1 \otimes M_2 N_2$ we end up with 4 contributions that don't cancel each other and can be rearranged into the commutator term $[A, A'] \otimes [B, B']$.

2) From the identity we have

$$\begin{aligned} \langle \Psi | \mathcal{B}^2 | \Psi \rangle &= 4 - \langle \Psi | [A, A'] \otimes [B, B'] | \Psi \rangle \\ &\leq 4 + |\langle \Psi | [A, A'] \otimes [B, B'] | \Psi \rangle| \end{aligned}$$

Now we must prove that the last term on the right-hand-side is less or equal to 4. Using the hint and Cauchy-Schwarz

$$|\langle \Psi | [A, A'] \otimes [B, B'] | \Psi \rangle| \leq \| [A, A'] \otimes I_B | \Psi \rangle \| \| I_A \otimes [B, B'] | \Psi \rangle \|$$

To estimate each norm on the right-hand-side we first use the triangle inequality

$$\| [A, A'] \otimes I_B | \Psi \rangle \| \leq \| AA' \otimes I_B | \Psi \rangle \| + \| A'A \otimes I_B | \Psi \rangle \|$$

and then

$$\begin{aligned} \| AA' \otimes I_B | \Psi \rangle \|^2 &= \langle \Psi | (AA' \otimes I_B)^\dagger (AA' \otimes I_B) | \Psi \rangle \\ &= \langle \Psi | (A'A \otimes I_B) (AA' \otimes I_B) | \Psi \rangle \\ &= \langle \Psi | A'A^2 A' \otimes I_B^2 | \Psi \rangle \\ &= \langle \Psi | A'^2 \otimes I_B | \Psi \rangle \\ &= \langle \Psi | I_A \otimes I_B | \Psi \rangle = 1 \end{aligned} \tag{1}$$

From which we deduce that the right-hand-side of the last inequality above equals $1 + 1 = 2$. Thus

$$\| [A, A'] \otimes I_B | \Psi \rangle \| \leq 2$$

and

$$|\langle \Psi | [A, A'] \otimes [B, B'] | \Psi \rangle| \leq 4$$

which completes the proof.

3) This last result can be justified from Cauchy-Schwarz:

$$\langle \Psi | \mathcal{B} | \Psi \rangle^2 \leq \| | \Psi \rangle \|^2 \| \mathcal{B} | \Psi \rangle \|^2 = \langle \Psi | \mathcal{B}^2 | \Psi \rangle$$

Since we proved before that the r.h.s is less or equal to 8 we get Tsirelson's bound for any state

$$\langle \Psi | \mathcal{B} | \Psi \rangle \leq 2\sqrt{2}$$

Note that the Bell state saturates the bound.